

Applications of Liouville Discrepancy to Quantum Dynamics

Matthew Bradshaw, Titus de Jong, Audrey Wang

REU in Semi-Algebraic Geometry in Schrödinger Equations.

July 17, 2025

Section 1

1 Quantum Dynamics

- Transport
- Quasi-Periodic Potentials

2 Discrepancy Estimates of Kronecker Sequences

- Examples of Discrepancy
- Bounds

3 Semi-Algebraic Sets

4 Application to Quantum Dynamics

The Schrödinger Equation

The Schrödinger Equation

- The dynamics of a quantum particle are described by a complex-valued *wavefunction*, ψ

The Schrödinger Equation

- The dynamics of a quantum particle are described by a complex-valued *wavefunction*, ψ
- Time evolution is governed by the *Schrödinger equation*

$$i\frac{\partial\psi}{\partial t} = \frac{\partial^2\psi}{\partial x^2} + V(x)\psi(x)$$

The Schrödinger Equation

- The dynamics of a quantum particle are described by a complex-valued *wavefunction*, ψ
- Time evolution is governed by the *Schrödinger equation*

$$i\frac{\partial\psi}{\partial t} = \frac{\partial^2\psi}{\partial x^2} + V(x)\psi(x)$$

- This can be spatially discretized (eg. for a crystal lattice, numerical simulation), to obtain

$$i\frac{\partial\psi_n}{\partial t} = (\psi_{n+1} - 2\psi_n + \psi_{n-1}) + V_n\psi_n$$

where now ψ, V are sequences

Quantum Dynamics - Schrödinger Operators

Definition 1.1: Schrödinger operator

A **Schrödinger operator** $H : \ell^2(\mathbb{Z}) \rightarrow \ell^2(\mathbb{Z})$ acts on a *wavefunction* $\psi : \mathbb{Z} \rightarrow \mathbb{C}$ according to

$$(H\psi)_n = \psi_{n+1} + \psi_{n-1} + V_n\psi_n, \quad (1)$$

where $V : \mathbb{Z} \rightarrow \mathbb{R}$ is the *potential*

Definition 1.1: Schrödinger operator

A **Schrödinger operator** $H : \ell^2(\mathbb{Z}) \rightarrow \ell^2(\mathbb{Z})$ acts on a *wavefunction* $\psi : \mathbb{Z} \rightarrow \mathbb{C}$ according to

$$(H\psi)_n = \psi_{n+1} + \psi_{n-1} + V_n\psi_n, \quad (1)$$

where $V : \mathbb{Z} \rightarrow \mathbb{R}$ is the *potential*

- Comes from discretizing the 1D Hamiltonian operator: $\frac{\partial^2}{\partial x^2} + V(x)$.
- Our goal is to study $\psi(n, t)$ given by

$$i\frac{\partial\psi}{\partial t} = H\psi. \quad (2)$$

Definition 1.1: Schrödinger operator

A **Schrödinger operator** $H : \ell^2(\mathbb{Z}) \rightarrow \ell^2(\mathbb{Z})$ acts on a *wavefunction* $\psi : \mathbb{Z} \rightarrow \mathbb{C}$ according to

$$(H\psi)_n = \psi_{n+1} + \psi_{n-1} + V_n\psi_n, \quad (1)$$

where $V : \mathbb{Z} \rightarrow \mathbb{R}$ is the *potential*

- Comes from discretizing the 1D Hamiltonian operator: $\frac{\partial^2}{\partial x^2} + V(x)$.
- Our goal is to study $\psi(n, t)$ given by

$$i\frac{\partial\psi}{\partial t} = H\psi. \quad (2)$$

- Given $\psi(0) = \phi$, this has the exact solution $\psi(t) = e^{-itH}\phi$.

Definition 1.1: Schrödinger operator

A **Schrödinger operator** $H : \ell^2(\mathbb{Z}) \rightarrow \ell^2(\mathbb{Z})$ acts on a *wavefunction* $\psi : \mathbb{Z} \rightarrow \mathbb{C}$ according to

$$(H\psi)_n = \psi_{n+1} + \psi_{n-1} + V_n\psi_n, \quad (1)$$

where $V : \mathbb{Z} \rightarrow \mathbb{R}$ is the *potential*

- Comes from discretizing the 1D Hamiltonian operator: $\frac{\partial^2}{\partial x^2} + V(x)$.
- Our goal is to study $\psi(n, t)$ given by

$$i\frac{\partial\psi}{\partial t} = H\psi. \quad (2)$$

- Given $\psi(0) = \phi$, this has the exact solution $\psi(t) = e^{-itH}\phi$.
- $|\psi_n|^2$ represents the probability of measurement at site n .

Quantum Dynamics - Schrödinger Operators

Quantum Dynamics - Schrödinger Operators

A Schrödinger operator can be expressed as an infinite matrix

$$H = \begin{pmatrix} \ddots & & & & & \\ & V_{-2} & 1 & 0 & 0 & 0 \\ & 1 & V_{-1} & 1 & 0 & 0 \\ & 0 & 1 & V_0 & 1 & 0 \\ & 0 & 0 & 1 & V_1 & 1 \\ & 0 & 0 & 0 & 1 & V_2 \\ & & & & & \ddots \end{pmatrix}$$

Quantum Dynamics - Schrödinger Operators

A Schrödinger operator can be expressed as an infinite matrix

$$H = \begin{pmatrix} \ddots & & & & & \\ & V_{-2} & 1 & 0 & 0 & 0 \\ & 1 & V_{-1} & 1 & 0 & 0 \\ & 0 & 1 & V_0 & 1 & 0 \\ & 0 & 0 & 1 & V_1 & 1 \\ & 0 & 0 & 0 & 1 & V_2 \\ & & & & & \ddots \end{pmatrix}$$

We can consider its restriction to a finite interval, eg. $[0,3]$:

$$H_{[0,3]} = \begin{pmatrix} V_0 & 1 & 0 & 0 \\ 1 & V_1 & 1 & 0 \\ 0 & 1 & V_2 & 1 \\ 0 & 0 & 1 & V_3 \end{pmatrix}$$

Quantum Dynamics - Example 1

$$\phi_n = e^{-n^2/2} \quad V_n = 0$$

Quantum Dynamics - Example 1

$$\phi_n = e^{-n^2/2} \quad V_n = 0$$

Insert diffusion example

Quantum Dynamics - Example 2

$$\phi_n = e^{ni-(n+\mu)^2/2} + e^{-ni-(n-\mu)^2/2} \quad V_n = 0$$

Quantum Dynamics - Example 2

$$\phi_n = e^{ni-(n+\mu)^2/2} + e^{-ni-(n-\mu)^2/2} \quad V_n = 0$$

Insert interference example

Quantum Dynamics - Example 3

$$\phi_n = e^{ni - (n+\mu)^2/2} \quad V_n = \begin{cases} 1, & 140 \leq n \leq 150 \\ 0, & \text{otherwise} \end{cases}$$

Quantum Dynamics - Example 3

$$\phi_n = e^{ni - (n+\mu)^2/2} \quad V_n = \begin{cases} 1, & 140 \leq n \leq 150 \\ 0, & \text{otherwise} \end{cases}$$

Insert reflection example

Definition 1.2: p^{th} moment

For an initial state ϕ , we define the p^{th} **moment** as:

$$\langle |X_H|^p | \phi \rangle(t) = \sum_{n \in \mathbb{Z}} |n|^p |\psi_n(\tau)|^2$$

Then, we define the **average p^{th} moment** as:

$$\langle |\tilde{X}_H|^p | \phi \rangle(t) = \frac{2}{t} \int_0^\infty e^{2\tau/t} \sum_{n \in \mathbb{Z}} |n|^p |\psi_n(\tau)|^2 d\tau$$

- Captures how $\psi(t) = e^{-itH}\phi$ spreads out over time

Transport - Example 1

If V is periodic, the p^{th} moment grows like

$$\langle X_H |_{\phi}^p \rangle \sim t^p$$

This is known as *ballistic transport* [Fil20].

Transport - Example 1

If V is periodic, the p^{th} moment grows like

$$\langle X_H |_{\phi}^p \rangle \sim t^p$$

This is known as *ballistic transport* [Fil20].

Insert periodic transport example

Transport - Example 2

If V is random, the p^{th} moment is bounded:

$$\langle X_H |_{\phi}^p \rangle \leq C \quad \text{for some } C > 0$$

This is known as *dynamic localization* [Klein].

Insert random transport example

Penrose Tilings

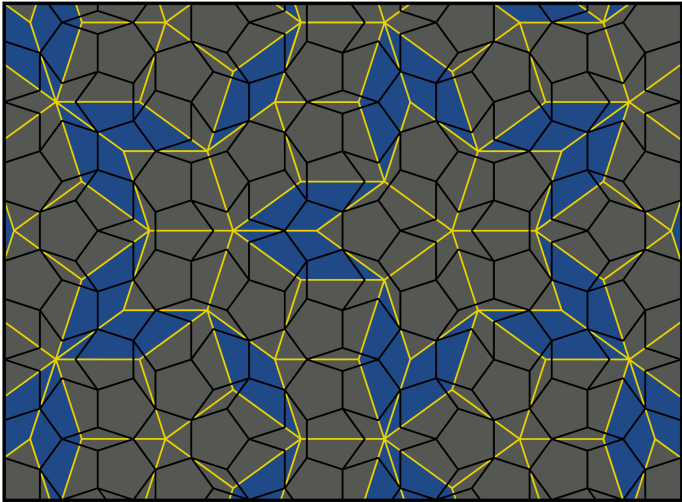


Figure: Image By Inductiveload - Own work, Public Domain, [Link](#)

The Fibonacci Hamiltonian

Consider a sequence S_N of words where $S_0 = 0$ and S_N is defined as the N -th application of the substitution rule

$$0 \mapsto 01 \quad \text{and} \quad 1 \mapsto 0$$

Taking the limit as $N \rightarrow \infty$, one obtains the Fibonacci Word (Hamiltonian) Eg.

$$0100101001001 \dots$$

Well studied model of Quasi-Crystals.

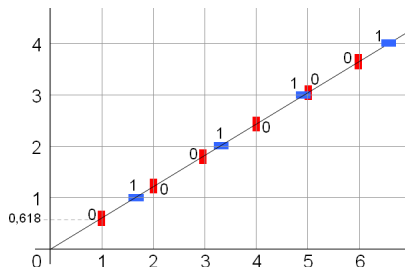


Figure: Line of Slope $\frac{1}{\phi}$. By Prokofiev - Own work, CC BY-SA 3.0, [Link](#)

Definition 1.3: Quasi-periodic

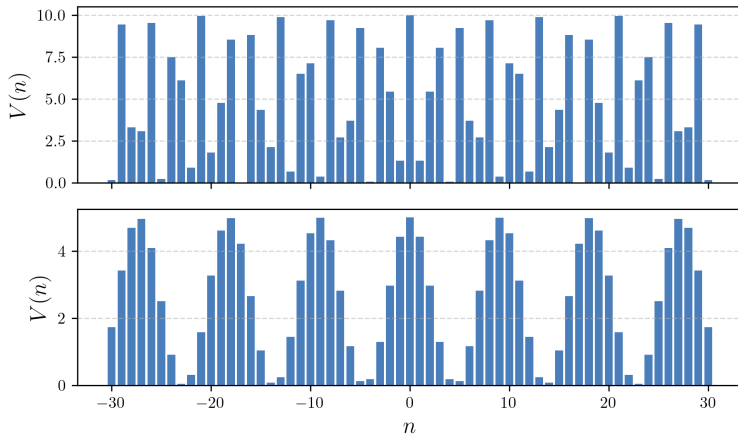
If $f : \mathbb{R} \rightarrow \mathbb{R}$ is a periodic, analytic function, and if α is an irrational number, we call

$$V(n) = f(n\alpha)$$

a **quasi-periodic** potential.

Quasi-Periodic Potentials

$$V(n) = \cos^2(n\pi\alpha)$$



$$\alpha = \frac{\sqrt{5}-1}{2}$$

$$\alpha = \sum_{k=1}^{\infty} \frac{1}{10^k!}$$

Continued Fractions

Fact 1.5: Contined fraction expansion

Every positive real number can be expressed as a continued fraction of the form

$$\alpha = a_0 + \frac{1}{a_1 + \frac{1}{a_2 + \dots}}$$

for $a_i \in \mathbb{N}$, and we write $\alpha = [a_0; a_1, a_2, \dots]$.

Fact 1.6: Contined fraction expansion

Every positive real number can be expressed as a continued fraction of the form

$$\alpha = a_0 + \frac{1}{a_1 + \frac{1}{a_2 + \dots}}$$

for $a_i \in \mathbb{N}$, and we write $\alpha = [a_0; a_1, a_2, \dots]$.

- Rational α have finite expansions

Fact 1.7: Contined fraction expansion

Every positive real number can be expressed as a continued fraction of the form

$$\alpha = a_0 + \frac{1}{a_1 + \frac{1}{a_2 + \dots}}$$

for $a_i \in \mathbb{N}$, and we write $\alpha = [a_0; a_1, a_2, \dots]$.

- Rational α have finite expansions
- Irrational α have infinite expansions

Fact 1.8: Contined fraction expansion

Every positive real number can be expressed as a continued fraction of the form

$$\alpha = a_0 + \frac{1}{a_1 + \frac{1}{a_2 + \dots}}$$

for $a_i \in \mathbb{N}$, and we write $\alpha = [a_0; a_1, a_2, \dots]$.

- Rational α have finite expansions
- Irrational α have infinite expansions
 - Terminating expansion gives the best rational approximation

Fact 1.9: Contined fraction expansion

Every positive real number can be expressed as a continued fraction of the form

$$\alpha = a_0 + \frac{1}{a_1 + \frac{1}{a_2 + \dots}}$$

for $a_i \in \mathbb{N}$, and we write $\alpha = [a_0; a_1, a_2, \dots]$.

- Rational α have finite expansions
- Irrational α have infinite expansions
 - Terminating expansion gives the best rational approximation
 - Growth rate of a_i determines how closely approximated by rationals α is.

Fact 1.10: Contined fraction expansion

Every positive real number can be expressed as a continued fraction of the form

$$\alpha = a_0 + \frac{1}{a_1 + \frac{1}{a_2 + \dots}}$$

for $a_i \in \mathbb{N}$, and we write $\alpha = [a_0; a_1, a_2, \dots]$.

- Rational α have finite expansions
- Irrational α have infinite expansions
 - Terminating expansion gives the best rational approximation
 - Growth rate of a_i determines how closely approximated by rationals α is.

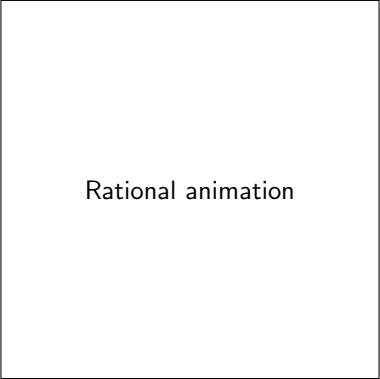
$$\frac{\sqrt{5}+1}{2} = [1; 1, 1, 1, \dots] \quad \sum_{k=1}^{\infty} \frac{1}{10^{k!}} = [0; 9, 11, 99, 1, 10, 9, 999999999999, 1, \dots]$$

Diophantine Conditions

Consider what happens when we plot the sequence $e^{2\pi i n \alpha}$ for rational α :

Diophantine Conditions

Consider what happens when we plot the sequence $e^{2\pi i n \alpha}$ for rational α :



Rational animation

$$\alpha = \frac{3}{11}$$

Diophantine Conditions

$$x_n = e^{2\pi i n \alpha}$$

ϕ animation

Liouville animation

$$\alpha = \frac{\sqrt{5} + 1}{2}$$

$$\alpha = \sum_{n=1}^{\infty} \frac{1}{10^{k!}}$$

Diophantine Conditions

Diophantine Conditions

Definition 1.7: Diophantine

We call $\alpha \in \mathbb{R}$ **diophantine** if there is some $\gamma > 0, \tau \geq 1$ such that for all $m \in \mathbb{Z}$,

$$\|m\alpha\| := \text{dist}(m\alpha, \mathbb{Z}) \geq \frac{\gamma}{|m|^\tau}$$

We call these *poorly approximable by rationals*, or colloquially, “more irrational”

Diophantine Conditions

Definition 1.7: Diophantine

We call $\alpha \in \mathbb{R}$ **diophantine** if there is some $\gamma > 0, \tau \geq 1$ such that for all $m \in \mathbb{Z}$,

$$\|m\alpha\| := \text{dist}(m\alpha, \mathbb{Z}) \geq \frac{\gamma}{|m|^\tau}$$

We call these *poorly approximable by rationals*, or colloquially, “more irrational”

Definition 1.8: Liouville

We call $m \in \mathbb{R}$ **Liouville** if, for all $\gamma > 0, \tau > 1$, there are infinitely many $m \in \mathbb{Z}$ for which

$$\|m\alpha\| < \frac{\gamma}{|m|^\tau}$$

We call these *too well approximable by rationals*, or colloquially, “less irrational”

Transport - Example 3

If α is diophantine, the transport is similar to the random case [BG00]

Insert diophantine transport example

Transport - Example 4

If α is Liouville, the transport is similar to the periodic case

Insert Liouville transport example

Main Results

We determine bounds on transport for a number of Diophantine-like conditions

Main Results

We determine bounds on transport for a number of Diophantine-like conditions

$$\|m\alpha\| \geq \frac{1}{|m|^\gamma} \quad \implies \quad \langle X_H|_\phi^p \rangle \leq (\log t)^{pc}$$

Main Results

We determine bounds on transport for a number of Diophantine-like conditions

$$\|m\alpha\| \geq \frac{1}{|m|^\gamma} \quad \implies \quad \langle X_H|_\phi^p \rangle \leq (\log t)^{pc}$$

$$\|m\alpha\| \geq e^{-\kappa(\log |n|)^\gamma} \quad \implies \quad \langle X_H|_\phi^p \rangle \leq e^{p(\log \log t)^c}$$

Main Results

We determine bounds on transport for a number of Diophantine-like conditions

$$\|m\alpha\| \geq \frac{1}{|m|^\gamma} \implies \langle X_H|_\phi^p \rangle \leq (\log t)^{pc}$$

$$\|m\alpha\| \geq e^{-\kappa(\log |n|)^\gamma} \implies \langle X_H|_\phi^p \rangle \leq e^{p(\log \log t)^c}$$

$$\|m\alpha\| \geq e^{-\kappa|n|^\gamma} \implies \langle X_H|_\phi^p \rangle \leq e^{p(\log t)^c}$$

Section 2

1 Quantum Dynamics

- Transport
- Quasi-Periodic Potentials

2 Discrepancy Estimates of Kronecker Sequences

- Examples of Discrepancy
- Bounds

3 Semi-Algebraic Sets

4 Application to Quantum Dynamics

Definition 2.1: Discrepancy

Let $\{\mathbf{x}_n\}_{n \in \mathbb{N}}$ be a sequence in $[0, 1)^d$. Define $P_N = \{\mathbf{x}_1, \mathbf{x}_2, \dots, \mathbf{x}_N\}$ then define the **Discrepancy** of P_N as

$$D_N(P_N) = \sup_{I \subset [0,1)^d} \left| \frac{1}{N} \sum_{n=1}^N \chi_I(x_n) - \text{Vol}(I) \right| \quad (3)$$

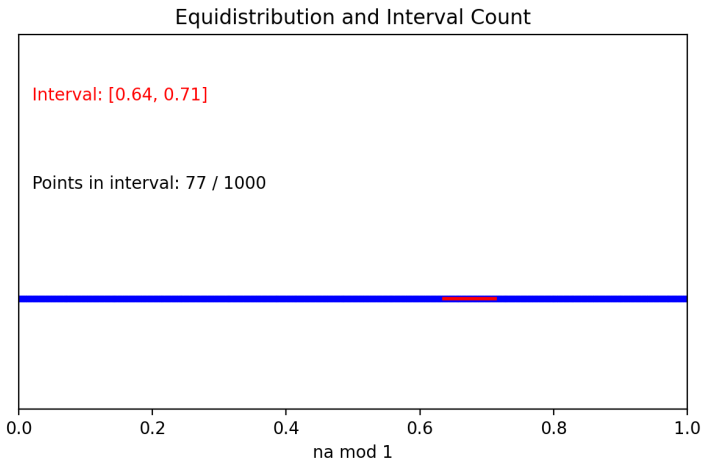
Where I is an axis-aligned rectangle.

Definition 2.2: Kronecker Sequence

Let $\alpha = (\alpha_1, \alpha_2, \dots, \alpha_d) \in \mathbb{R}^d$. The **Kronecker Sequence** $\{\mathbf{x}_n\}_{n \in \mathbb{N}}$ of α is defined via the relation:

$$\mathbf{x}_n := n\alpha \pmod{\mathbb{Z}^d} \quad (4)$$

Example 1 in \mathbb{R}

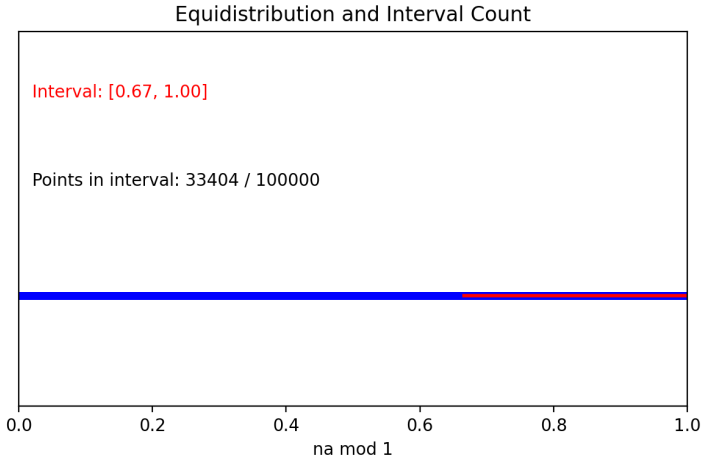


$$a = \sqrt{2}$$

$$\text{Expected: } (0.71 - 0.64) \cdot 1000 = 70$$

$$\text{Discrepancy: } \frac{77-70}{1000} = 0.007$$

Example 2 in \mathbb{R}

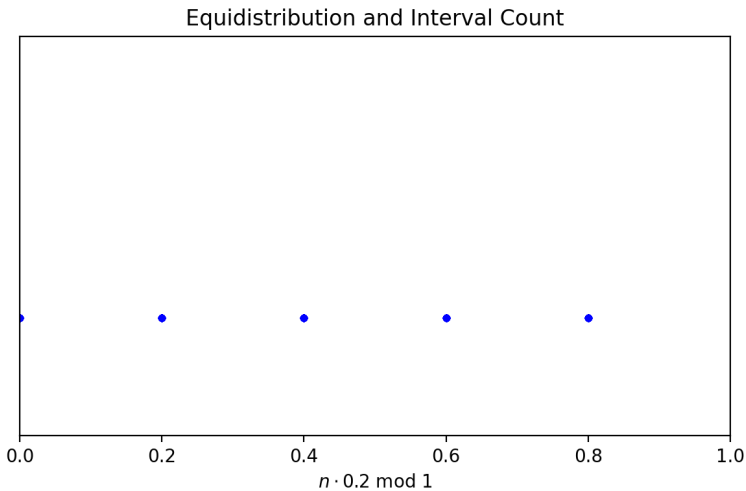


$$a = \sqrt{2}$$

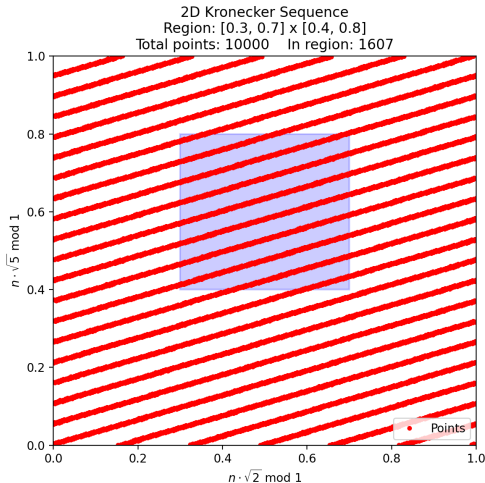
$$\text{Expected: } (1 - 0.67) \cdot 100000 = 33000$$

$$\text{Discrepancy: } \frac{33404 - 33000}{100000} = 0.00404$$

Example 3 in \mathbb{R}



Example 4 in \mathbb{R}^2

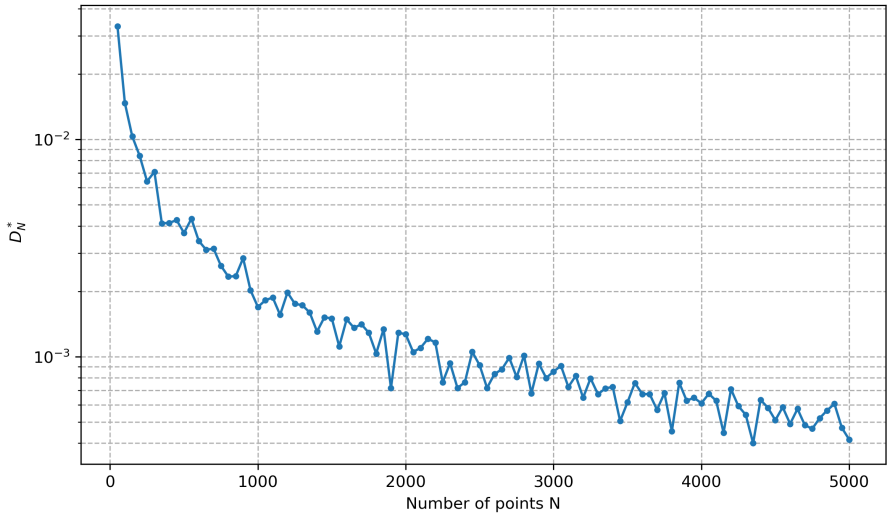


Expected: $((0.7 - 0.3) \cdot (0.8 - 0.4)) \cdot 10000 = 1600$

Discrepancy: $\frac{1607 - 1600}{10000} = 0.0007$

Discrepancy as $N \rightarrow \infty$

Discrepancy of Kronecker sequence $(\{n\sqrt{2}\})$ as $N \rightarrow \infty$



Theorem 2.3: Denseness of Kronecker Sequences

Let $\alpha = (\alpha_1, \alpha_2, \dots, \alpha_d) \in \mathbb{R}^d$. Then the Kronecker sequence $\{n\alpha\}_{n \in \mathbb{N}}$ is dense in the unit torus $[0, 1)^d$ iff $\{1, \alpha_1, \dots, \alpha_d\}$ is linearly independent over the integers. We henceforth call this condition **rationally independent**.

- Hence, $D_N(\{n\alpha\}) \rightarrow 0$ as $N \rightarrow \infty$ if and only if α is rationally independent.
- Interested in leading order approximation of D_N (logarithmic, exponential, polynomial?)
- Heavily depends on whether α is Diophantine or Liouville.

Theorem ([DT97])

Suppose $\alpha \in \mathbb{R}^d$ satisfies

$$\|\mathbf{m}\alpha\| = \text{dist}(\mathbf{m} \cdot \alpha, \mathbb{Z}) \gtrsim \frac{\gamma}{|\mathbf{m}|^a} \quad (5)$$

for all $\mathbf{m} \in \mathbb{Z}^d$ for some $a \geq 1, \gamma > 0$. Then for some $\delta, C > 0$

$$D_N(\{\mathbf{x}_n\}) \leq CN^{-\delta} \quad (6)$$

Theorem

Suppose $\alpha \in \mathbb{R}^d$ is such that for some $\gamma, C > 0$

$$\|\mathbf{m}\alpha\| \gtrsim e^{-C|\mathbf{m}|^\gamma} \quad (7)$$

for all $\mathbf{m} \in \mathbb{Z}^d$ and for some $\gamma, C > 0$. Then for some $C_1 > 0$

$$D_N(\{\mathbf{x}_n\}) \leq C_1(\log N)^{-1/\gamma} \quad (8)$$

Proof of Theorem

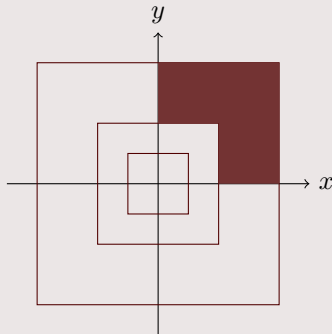
Proof.

- 1 Partition \mathbb{Z}^d as in figure.
- 2 Bound each partition, and count the number within each partition.
- 3 Apply to Erdős-Turan inequality
- 4 Summing values together yields:

$$D_N \lesssim \frac{1}{M} + \frac{1}{N} + \frac{e^{(C(M-1)^{d\gamma})}(M-1)^d}{2N}$$

- 5 Choose $M = e^{5CN^\gamma}$. Yields

$$D_N \leq C_1(\log N)^{-1/\gamma}$$



Theorem

Suppose $\alpha \in \mathbb{R}^d$ is such that for some $\gamma, C > 0$

$$\|\mathbf{m}\alpha\| \gtrsim e^{-C(\log |\mathbf{m}|)^\gamma} \quad (9)$$

for all $\mathbf{m} \in \mathbb{Z}^d$. Then for some $\delta, C > 0$

$$D_N(\{\mathbf{x}_n\}) \leq \begin{cases} Ce^{-(\log N)^{1/\gamma}} & \gamma \geq 1 \\ CN^{-\delta} & \gamma < 1 \end{cases} \quad (10)$$

Counting using Discrepancy

Recall

$$D_N(P_N) = \sup_{I \subset [0,1)^d} \left| \frac{1}{N} \sum_{n=1}^N \chi_I(n\alpha) - \text{Vol}(I) \right| \quad (11)$$

Hence, for a given rectangle

$$\sum_{n=1}^N \chi_I(n\alpha) \leq N(D_N - \text{Vol}(I)) \quad (12)$$

Section 3

1 Quantum Dynamics

- Transport
- Quasi-Periodic Potentials

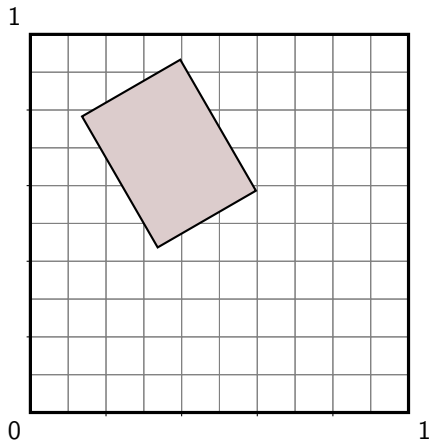
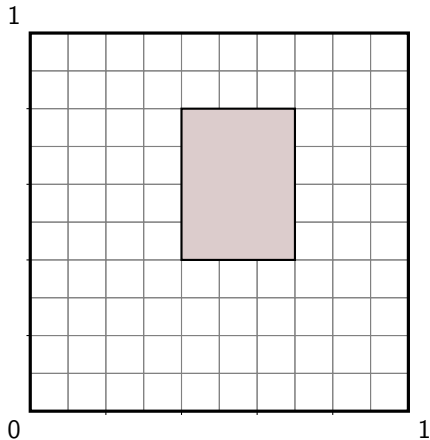
2 Discrepancy Estimates of Kronecker Sequences

- Examples of Discrepancy
- Bounds

3 Semi-Algebraic Sets

4 Application to Quantum Dynamics

Expansion to Semi-Algebraic Sets



Semi-Algebraic Sets

Smiley Face with Semi-Algebraic Regions

Definition 3.1: Semi-Algebraic Sets

A set \mathcal{S} is **semi-algebraic** if it is a finite union of polynomial inequalities and equalities.



Head: $x^2 + y^2 \leq 1$

Left eye: $(x + 0.4)^2 + (y - 0.4)^2 \leq 0.01$

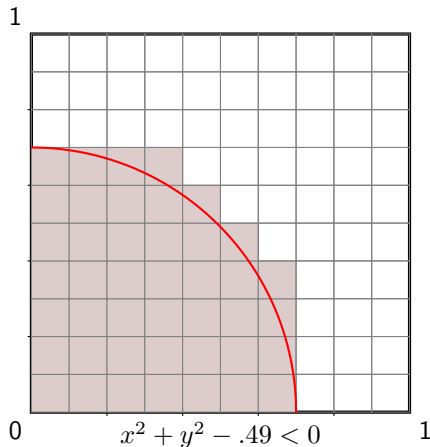
Right eye: $(x - 0.4)^2 + (y - 0.4)^2 \leq 0.01$

Smile: $1.5x^2 - 0.7 < y < 1.5x^2 - 0.6, \quad |x| < 0.5$

Expansion to Semi-Algebraic Sets

Theorem 3.2: [Bou05]

Let $\mathcal{S} \subseteq [0,1]^n$ be a semi-algebraic set of degree B . Let $\text{Vol}(\mathcal{S}) < \epsilon^n$ for some $\epsilon > 0$. \mathcal{S} can be covered by $B^C(\epsilon)^{1-n}$ ϵ -balls.

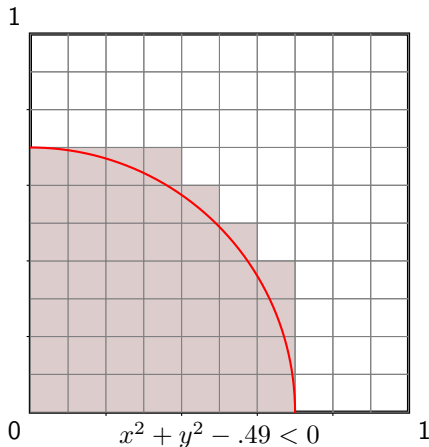


Expansion to Semi-Algebraic Sets

Theorem 3.3: [Bou05]

Let $\mathcal{S} \subseteq [0,1]^n$ be a semi-algebraic set of degree B . Let $\text{Vol}(\mathcal{S}) < \epsilon^n$ for some $\epsilon > 0$. \mathcal{S} can be covered by $B^C(\epsilon)^{1-n}$ ϵ -balls.

- Recall: V is analytic, hence V is well approximated by polynomials



Section 4

1 Quantum Dynamics

- Transport
- Quasi-Periodic Potentials

2 Discrepancy Estimates of Kronecker Sequences

- Examples of Discrepancy
- Bounds

3 Semi-Algebraic Sets

4 Application to Quantum Dynamics

Definition 4.1: Eigenvalues

Let $A \in M_{n \times n}(\mathbb{R})$. An eigenvalue- eigenvector pair are $\lambda \in \mathbb{C}$, $x \in \mathbb{R}^n$ which satisfy:

$$Ax = \lambda x$$

Alternatively, $\lambda \in \mathbb{C}$ is an eigenvalue iff

$$(A - \lambda I_n)^{-1} \text{ does not exist}$$

Definition 4.2: Eigenvalues

Let $A \in M_{n \times n}(\mathbb{R})$. An eigenvalue- eigenvector pair are $\lambda \in \mathbb{C}$, $x \in \mathbb{R}^n$ which satisfy:

$$Ax = \lambda x$$

Alternatively, $\lambda \in \mathbb{C}$ is an eigenvalue iff

$$(A - \lambda I_n)^{-1} \text{ does not exist}$$

- $H\psi = E\psi$
- Approximate H with finite dimensional H_N .

Connection to Quantum Dynamics

- Define $G_N(z; \theta) := (H_N(\theta) - zI_N)^{-1}$. Called the **Green's Function** of H .
- Since V is analytic, approximate with high-degree polynomial.
- Points where G_N doesn't exist are contained in Semi-Algebraic Set (Denote Θ_N).

Theorem 4.3: [HZ22] Large Deviation Estimates

$$\text{Leb} \left\{ \theta \in [0, 1]^d : \det(H_N(\theta) - zI) < e^{-N^{1/2}} \right\} \leq e^{-N^{1/4}}$$

Theorem 4.4: Counting Estimates

$$\#\{|n| \leq N : \theta + n\alpha \in \Theta_N\} \text{ is } o(N^{-1})$$

Theorem 4.5

If α satisfies

$$\|m\alpha\| > e^{-(\log |m|)^\gamma} \quad (13)$$

Then,

$$\langle |X|_\phi^p \rangle(t) \lesssim e^{p(\log \log t)^c} \quad (14)$$

Theorem 4.6

If α satisfies

$$\|m\alpha\| > e^{-|m|^\gamma} \quad (15)$$

Then,

$$\langle |X|_\phi^p \rangle(t) \lesssim e^{p(\log t)^c} \quad (16)$$

Proof of Theorem

Proof.

Lemma 4.7

Let $z = E + \frac{i}{t}$. Assume that α is Liouvillian. Then for any j with $|j| \leq M$,

$$|G(j, n)| \lesssim t^4 e^{-\frac{c}{20}|I|}$$

- Define

$$a(j, n, t) = \frac{2}{t} \int_0^\infty e^{2\tau/t} |(e^{-i\tau H} \delta_j, \delta_n)|^2 d\tau$$

- Hence,

$$\langle |X|_\phi^p \rangle = \sum_{n \in \mathbb{Z}} \sum_{|j| \leq M} |n|^p a(j, n, t)$$



Proof of Theorem

Proof.

- By Parseval's Identity,

$$a(j, n, t) = \frac{1}{\pi t} \int_{\mathbb{R}} |G(z)(j, n)|^2 dz.$$

- Combes-Thomas Estimate, and the previous Lemma, we may estimate

$$a(j, n, t) \leq t^7 e^{-2c(\log(N))^C}$$

- By definition

$$\langle |X_H|_{\phi}^p \rangle(t) = \left(\sum_{|n| \leq R} + \sum_{|n| > R} \right) \sum_{|j| \leq M} |n|^p a(j, n, t)$$

- By construction

$$\langle |X_H|_{\phi}^p \rangle(t) \lesssim R^p + \sum_{|n| > R} \sum_{|j| \leq M} |n|^p a(j, n, t)$$



Proof of Theorem

Proof.

- Choose $R = e^{(\log T)^c}$
- Hence,

$$\langle |X_H|_\phi^p \rangle(t) \lesssim e^{p(\log T)^c}$$



- Using this methodology, we may extend to general case (up to choice of R).

Conclusion

- 1 Quasi-Periodic Schrödinger Operators are fundamentally unique.
- 2 Summary of Tools used to find upper bounds.
 - 1 Discrepancy
 - 2 Semi-Algebraic Geometry
 - 3 Spectral Theory
- 3 Derived upper bounds show range of growth rates.

Open Problems

- 1 Lower Bounds
- 2 Does there exist an α such that $\langle |X|_\phi^p \rangle \sim T^{p/2}$?

Acknowledgments

This work is carried out at the *Texas A&M Mathematics REU 2025* under the thematic program *Semi-Algebraic Geometry in Schrödinger Equations* funded by the NSF REU site grant DMS-2150094.

Part of this work was completed during the *2025 Texas Summer School in Mathematical Physics* sponsored by the NSF grants DMS-2052572, DMS-2246031, and DMS-2513006. The presenters thank Jake Fillman, Wencai Liu and Xin Liu for organizing the event.

The presenters thank Wencai Liu, Xueyin Wang, and Bingheng Yang for their mentorship throughout the REU.

Bibliography I

- [DT97] M. Drmota and R.F. Tichy. *Sequences, discrepancies and applications*. Vol. 1651. Lecture Notes in Mathematics. Springer-Verlag, Berlin, 1997, pp. xiv+503. ISBN: 3-540-62606-9. DOI: 10.1007/BFb0093404. URL: <https://doi.org/10.1007/BFb0093404>.
- [BG00] Jean Bourgain and Michael Goldstein. *On nonperturbative localization with quasi-periodic potential*. 2000. arXiv: math-ph/0011053 [math-ph]. URL: <https://arxiv.org/abs/math-ph/0011053>.
- [Bou05] J. Bourgain. *Green's function estimates for lattice Schrödinger operators and applications*. Vol. 158. Annals of Mathematics Studies. Princeton University Press, Princeton, NJ, 2005, pp. x+173. ISBN: 0-691-12098-6. DOI: 10.1515/9781400837144. URL: <https://doi.org/10.1515/9781400837144>.
- [Fil20] Jake Fillman. *Ballistic Transport for Periodic Jacobi Operators on \mathbb{Z}^d* . 2020. arXiv: 2011.11522 [math.SP]. URL: <https://arxiv.org/abs/2011.11522>.

- [HZ22] R. Han and S. Zhang. “Large deviation estimates and Hölder regularity of the Lyapunov exponents for quasi-periodic Schrödinger cocycles”. In: *Int. Math. Res. Not. IMRN* 3 (2022), pp. 1666–1713. ISSN: 1073-7928, 1687-0247. DOI: 10.1093/imrn/rnz319. URL: <https://doi.org/10.1093/imrn/rnz319>.

Thank You!